

Self-Imaging Waves

Joint work w/ J. Garnier, L. Giovangigli, Q. Goepfert, E. Parolin



Probing the speckle to estimate the effective speed of sound, a first step towards quantitative ultrasound imaging



Institut Langevin ONDES ET IMAGES ESPCI E PARIS PSL

Adaptative optics in 1 min

Theoretical model for a linear imaging system

Theoretical Point Spread Function (PSF)

Adaptative optics in 1 min

Theoretical model for a linear imaging system

Practical situation with point-like object in field of view

Adaptative Optics

Theoretical Point Spread Function (PSF)

Fitting of the PSF

Adjust perturbative parameters (aberrations, travel time...) to find the focusing law

(Compensating modelisation approximations)





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(Compensating modelisation approximations)



High Resolution Imaging in Microscopy and Ophthalmology

Forewords by Stefan W. Hell and Robert N. Weinreb 🖄 Springer OPEN



Flavien Bureau: Analyse multi-dimensionnelle de la matrice de réflexion pour l'imagerie ultrasonore quantitative



Image is computed using space time correspondance





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In practice ?

Speed of sound in soft tissues varies ($\pm 10\,\%$) around a known value Soft tissues are « weakly scattering »





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Constant speed of sound Single diffusion regime





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Improvements

Access to S.O.S

SNR, contrast, resolution, diagnosis...







Signal is the backscattered field

Image is computed using space time correspondance



Speed of sound in soft tissues varies ($\pm 10\%$) around a known value Soft tissues are « weakly scattering »

Practical Hypotheses

Constant speed of sound Single diffusion regime

Improvements

Access to S.O.S

SNR, contrast, resolution, diagnosis...

No easy access to travel times







Backscattered signal is generated by discontinuities of acoustic impedance $Z=\sqrt{}$





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Discontinuities appear at a microscopic scale (hundreds/wavelength), no isolated reflectors





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Reflection Matrix Imaging for Wave Velocity Tomography

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(Dated: September 24, 2024)

New method to reconstruct the propagation speed of ultrasounds in soft tissues inspired by adaptative optics

Provide a robust mathematical model for this inverse problem

PSF probing in the speckle



Experimental work by A. Aubry and his team at Institut Langevin.







1. MODEL / DIRECT PROBLEM

2. INVERSE PROBLEM

Content



1. MODEL / DIRECT PROBLEM

2. INVERSE PROBLEM

Content

What is is that we measure

What do we do with those measurements

Content

1. Stochastic Homogenization of the Wave Equation

1.1.The Scattering Problem

- 1.2. Model for Soft Tissues
- 1.3. Homogenization & Representation Of The Scattered Field
- 2. Ultrasound Imaging Analysis
 - 2.1. Mathematical Model
 - 2.2. The Focal Spot
 - 2.3. Sound Speed Estimators
 - 2.4. Focusing in the speckle

The acoustic scattering problem



$$\begin{cases} \nabla \cdot (a\nabla u) + \omega^2 nu = 0\\ \lim_{|x| \to \infty} |x|^{\frac{d-1}{2}} \left(\partial_{|x|} \left(u - u^i \right) - ik_0 \left(u - u^i \right) \right) = 0 \end{cases}$$



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Physical parameters:

$$\begin{aligned} &(x) := a_0 \mathbb{1}_{\mathbb{R}^d \setminus \overline{D}} + a_D(x) \mathbb{1}_D, \\ &(x) := n_0 \mathbb{1}_{\mathbb{R}^d \setminus \overline{D}} + n_D(x) \mathbb{1}_D. \\ &= \omega \sqrt{n_0 a_0^{-1}} \end{aligned}$$



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What is a good model for a(x) and n(x)?



Soft tissues are incredibly complex

Realistic model for a(x) and n(x) is a lost cause

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Tissue-mimicking phantoms are much simpler :

LOTS of small scatterers densely embedded in an homogeneous matrix

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Variation in speed of sound obtained by locally changing the repartition of scatterers

Mathematical framework requirements

Random Stable w/ respect to size and number of scatterers Can infer macroscopic quantities from the micro-structure

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Stochastic Homogenization

 \mathbb{R}^{d} \



 $n u - u^i$

Mathematical description of the micro-structured medium

- \mathcal{Q}_i Domains of measure 1
- x_i Random centers distributed in \mathbb{R}^d

 $S = \bigcup_i \{x_i + Q\}$ Set of size 1 scatterers

$$\begin{cases} a_1 := a_M \mathbb{1}_{\mathbb{R}^d \setminus \overline{S}} + a_S \mathbb{1}_S, \\ n_1 := n_M \mathbb{1}_{\mathbb{R}^d \setminus \overline{S}} + n_S \mathbb{1}_S. \end{cases}$$

$$\forall x \in D \begin{cases} a_D^{\varepsilon}(x) = a_1\left(\frac{x}{\varepsilon}\right), \\ n_D^{\varepsilon}(x) := n_1\left(\frac{x}{\varepsilon}\right). \end{cases}$$





Example with $\varepsilon_1 > \varepsilon_2 > \varepsilon_3$

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Mathematical description of the micro-structured medium

 Q_i Domains of measure 1 x_i Random centers distributed in \mathbb{R}^d

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Example with $\varepsilon_1 > \varepsilon_2 > \varepsilon_3$







For the case $a \not\equiv 1$ see L. Giovangigli's talk

Simplification $a \equiv 1$





Governing equation total field u

$$\begin{cases} \Delta u + \omega^2 (n_0 + (n_{\varepsilon} - n_0) \mathbb{1}_D) u = 0\\ \lim_{|x| \to \infty} |x|^{\frac{d-1}{2}} \left(\partial_{|x|} \left(u - u^i \right) - i k_0 \left(u - u^i \right) \right) = 0. \end{cases}$$

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Governing equation coherent field u_0

$$\begin{cases} \Delta u_0 + \omega^2 (n_0 + (n^* - n_0) \mathbb{1}_D) u_0 = 0\\ \lim_{|x| \to \infty} |x|^{\frac{d-1}{2}} \left(\partial_{|x|} \left(u_0 - u^i \right) - ik_0 \left(u_0 - u^i \right) \right) = 0 \end{cases}$$

Effective index in D

$$n^{\star} \equiv \mathbb{E}[n_{\varepsilon}]$$



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Effective index in D

$$n^{\star} \equiv \mathbb{E}[n_{\varepsilon}]$$

Speed of sound outside of D

$$k_0 := \omega \sqrt{n_0} := \frac{\omega}{c_0}$$

$$k^* := \omega \sqrt{n^*} := \frac{\omega}{c^*}$$
Effective speed of sound in D



Governing equation total field u

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Homogenized Green function

$$\begin{cases} \Delta G^{\star}(\cdot, y) + \omega^{2}(n_{0} + (\mathbf{n}^{\star} - n_{0})\mathbb{1}_{D})G^{\star}(\cdot, y) = \delta_{y} \\ \lim_{|x| \to \infty} |x|^{\frac{d-1}{2}} \left(\partial_{|x|}G^{\star}(x, y) - ik_{0}G^{\star}(x, y)\right) = 0. \end{cases}$$

Governing equation coherent field u_0

$$\begin{cases} \Delta u_0 + \omega^2 (n_0 + (\mathbf{n}^* - n_0) \mathbb{1}_D) u_0 = 0\\ \lim_{|x| \to \infty} |x|^{\frac{d-1}{2}} \left(\partial_{|x|} \left(u_0 - u^i \right) - ik_0 \left(u_0 - u^i \right) \right) = 0 \end{cases}$$

Effective index in D

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Speed of sound outside of D

$$k_0 := \omega \sqrt{n_0} := \frac{\omega}{c_0}$$

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Effective speed of sound in D



The scattered field

$$\mathcal{U}^s := \omega^2 \int_D \left(n_{\varepsilon}(y) - n^* \right) u_0(y) G^*(\cdot, y) \mathrm{d}y.$$

Quantitative Homogenization Theorem (Garnier, Giovangigli, Goepfert, M.)

$$\|u(x) - u^i(x) - \mathcal{U}^s(x)\|_{L^2(\Omega)}$$

 C_ω indep. of ε and x . $p' = \frac{2}{1}$

 $\begin{aligned} & (2) \leq C_{\omega} \varepsilon^{\frac{d}{p'}} & \forall x \in \mathcal{P} \\ & \frac{+\beta}{+\beta} & \text{w/} \sup_{x \in D} \int_{D} |G^{\star}(x,y)|^{2+\beta} \, \mathrm{d}y < \infty \end{aligned}$

Illustrations

Example of domain



Illustrations



Example of domain



Total field u^{ε}

Illustrations



Example of domain



Total field u^{ε}

- 1.9e+00	
- 16	
1.4	
- 1.Z 1	
- 0.0	
- 0.0	
- 0.4	
- 0.2	
- 0	
0.2	
– -0.4	
0.6	
0.8	
1	
1.2	
- -1 <u>4</u>	
1.6	
1.9e+00	

Coherent field u_0


Illustrations









Total field u^{ε}



Coherent field u_0



Illustrations









Total field u^{ε}

u_0
- 3.4e-01
- 0.25
– 0.2
– 0.15
– 0.1
- 0.05
— O
0.05
0.1
0.15
0.2
0.25
3.1e-01

Coherent field u_0



 \mathcal{U}^{s} (First order approximation)



Conclusion of part 1

It is possible to derive effective properties for the ultrasound propagation from the statistical properties of the micro-structure

It is possible to approximate the backscattered field as a function of the micro-structure and the coherent field

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Simple model

$$\begin{cases} \Delta u + \frac{\omega^2}{\left(c^{\star}\right)^2} \left(1 + \left(n_{\varepsilon} - 1\right) \mathbb{1}_D\right) u = 0\\ \lim_{|x| \to \infty} |x|^{\frac{d-1}{2}} \left(\partial_{|x|} \left(u - u^i\right) - i\frac{\omega}{c^{\star}} \left(u - u^i\right)\right) = 0. \end{cases}$$

Another simplification for the medium:

Governing equation

Simple model

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Another simplification for the medium:

$$n_{\varepsilon}(x) = 1 + \tilde{n}\left(\frac{x}{\varepsilon}\right)$$
 with $\mathbb{E}\left[\tilde{n}(x)\right] = 0$,

Governing equation

Simple model

$$\begin{cases} \Delta u + \frac{\omega^2}{\left(c^{\star}\right)^2} \left(1 + \left(n_{\varepsilon} - 1\right) \mathbb{1}_D\right) u = 0\\ \lim_{|x| \to \infty} |x|^{\frac{d-1}{2}} \left(\partial_{|x|} \left(u - u^i\right) - i\frac{\omega}{c^{\star}} \left(u - u^i\right)\right) = 0. \end{cases}$$

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non-restrictive », just technical simplification to get :

$$G^{\star}(\cdot,z) = \Gamma^{\frac{\omega}{c^{\star}}}(\cdot,z)$$
 and $u^i = u_0.$

Governing equation

Match between homogenized and outer medium

Transducer array



Incident wave : $u^i(x_e, \cdot, \omega) := \Gamma^{\frac{\omega}{c^{\star}}}(x_e, \cdot)$

Simulation of US experiment

Physical domain D

Transducer array



Incident wave : $u^i(x_e, \cdot, \omega) := \Gamma^{\frac{\omega}{c^{\star}}}(x_e, \cdot)$

Simulation of US experiment

Physical domain D

Transducer array

Incident wave :
$$u^i(x_e,\cdot,\omega):=\Gamma$$

Measurement matrix:

Frequency $\omega \in \mathcal{B}$ (bandwidth). Effective speed of sound c^* . Probe $\mathcal{P} = [-\frac{a}{2}, \frac{a}{2}] \times 0$,

Emission point $x_e \in \mathcal{P}$, reception point $x_r \in \mathcal{P}$.

Simulation of US experiment

Physical domain D

$$\frac{\omega}{c^{\star}}(x_e,\cdot)$$

$$M_{e,r}(\omega) := u^s(x_e, x_r, \omega)$$



Simulation of US experiment

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Imaging functional

$$\overline{\omega})\Gamma^{\frac{\omega}{c^{\star}}}(x_e,z)\Gamma^{\frac{\omega}{c^{\star}}}(x_r,z)\,\mathrm{d}x_e\mathrm{d}x_r\mathrm{d}\omega$$



Computed Image Imaging domain D'





Simulation of US experiment

Physical domain D



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Computed Image Imaging domain D'

What if c^{\star} is not known ?



Characterization of the problem

Geometry

Active array of sources / sensor outside the medium

Reflection geometry

Large domain

Goal

Find the effective speed of sound *in situ* (geometrical component of the focusing law)

Prior

Stationarity and ergodicity of the fluctuations

Regime

Random fluctuations of the medium: short correlation length & order 1 amplitude

Dense distribution of scatterers (no sparsity)

Constraints

Low computational cost

The point-spread function (PSF)

First order approximation for the scattered fields

$$M(x_e, x_r, \omega) = \omega^2 \int_D (r) dr$$

 $n(y) - n^{\star} G^{\star}(x_r, y) G^{\star}(x_e, y) \mathrm{d}y$

The point-spread function (PSF)

First order approximation for the scattered fields

$$M(x_e, x_r, \omega) = \omega^2 \int_D \left(n(y) - n^* \right) G^*(x_r, y) G^*(x_e, y) \mathrm{d}y$$

$$I^{\star}(z) = \int_{D} \left(n(y) - n^{\star} \right) F^{c^{\star}}(z, y) \mathrm{d}y, \qquad z \in D'$$

Imaging functional as a convolution

The point-spread function (PSF)

First order approximation for the scattered fields

$$M(x_e, x_r, \omega) = \omega^2 \int_D \left(n(y) - n^* \right) G^*(x_r, y) G^*(x_e, y) \mathrm{d}y$$

Imaging functional as a convolution

$$I^{\star}(z) = \int_{D} \left(n(y) - n^{\star} \right) F^{c^{\star}}(z, y) \mathrm{d}y, \qquad z \in D'$$

$$F^{\star}(z,y) = \int_{\mathcal{B}} \omega^2 \left(\int_{\mathcal{P}} \Gamma^{\frac{\omega}{c^{\star}}}(z,x_r) \overline{G^{\star}(y,x_r)} \mathrm{d}x_r \right)^2 \mathrm{d}\omega \qquad z,y \in D' \times D.$$

Kernel of the operator (Point Spread Function)

$$M(x_e, x_r, \omega) = \omega^2 \int_D (r) dr$$

The case $c \neq c^{\star}$

 $n(y) - n^{\star}) G^{\star}(x_r, y) G^{\star}(x_e, y) \mathrm{d}y$

(Unchanged)

$$M(x_e, x_r, \omega) = \omega^2 \int_D \left(n(y) - n^* \right) G^*(x_r, y) G^*(x_e, y) \mathrm{d}y$$

Imaging functional as a convolution

$$I^{\boldsymbol{c}}(z) = \int_{D} \left(n(y) - n^{\star} \right) F^{\frac{\boldsymbol{c}}{\boldsymbol{c}^{\star}}}(z, y) \mathrm{d}y, \qquad z \in D'$$

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The case $c \neq c^{\star}$

(Unchanged)

 $D' \leftarrow \varphi_c(D)$

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$$F^{\frac{c}{c^{\star}}}(z,y) = \int_{\mathcal{B}} \omega^2 \left(\int_{\mathcal{P}} \Gamma^{\frac{\omega}{c}}(z,x_r) \overline{G^{\star}(y,x_r)} \mathrm{d}x_r \right)^2 \mathrm{d}\omega \qquad z,y \in \mathbf{D'} \times D.$$

The case $c \neq c^{\star}$

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Kernel of the operator (Point Spread Function)

$$D' \longrightarrow D'$$
$$z \longmapsto \left| F^{\frac{c}{c^{\star}}}(z, y_0) \right|$$

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 $y_0 = (0.01, 0.06) \in D$



Localisation is moved

$$D' \longrightarrow D'$$
$$z \longmapsto \left| F^{\frac{c}{c^{\star}}}(z, y_0) \right|$$



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Definition : Focal spot at threshold δ

Consider $0 < \delta \ll 1$ and $y_0 \in D$. There exists a domain $D'_{\delta}(y_0) \subset D'$ such that $\int_{D'\setminus\overline{D'_{\delta}(y_0)}} \left| F^{c^*}(z,y_0) \right| \mathrm{d}z < \delta \int_{D'} \left| F^{c^*}(z,y_0) \right| \mathrm{d}z \quad \text{and} \quad \frac{|D'_{\delta}(y_0)|}{|D'|} \ll \delta.$



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Definition : Focal spot at threshold δ

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0.8

0.6

0.4

0.2

0



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Asymptotic analysis of the phase
$F^{\frac{c}{c^{\star}}}(z,y) = \int_{\mathcal{B}} \omega^2 \left(\int_{\mathcal{D}} \Gamma^{\frac{\omega}{c}}(z,x_r) \overline{G^{\star}(y,x_r)} \mathrm{d}x_r \right)^2 \mathrm{d}\omega \qquad z,y \in D' \times D.$

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Focal spot

 $z, y \in D' \times D.$

$$\frac{|z - x_r|}{c} - \frac{|y - x_r|}{c^{\star}} = \mathcal{O}\left(\frac{1}{\omega}\right)$$

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Paraxial asymptotic regime ($\eta \ll 1$)

Bandwidth $\mathcal{B} := \frac{\mathcal{B}_0}{\eta}$ with $\mathcal{B}_0 := [\omega_0 - \frac{B}{2}, \omega_0 + \frac{B}{2}]$ Probe $\mathcal{P} := \eta^{\frac{1}{2}} \mathcal{P}_0$ with $\mathcal{P}_0 := [-\frac{a_0}{2}, \frac{a_0}{2}]^{d-1}$ Points $y := (y_\eta^{\perp}, y^{\parallel}) := (\eta^{\frac{1}{2}} y^{\perp}, y^{\parallel})$ Scatterers $\varepsilon := o(\eta)$ Focal spot

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$ z-x_r $	$ y-x_r $	$-\mathcal{O}\left(\frac{1}{-}\right)$
c	- $ -$	$- \left(\frac{-}{\omega} \right)$

Typical value
$[2, 6] \mathrm{MHz}$
$1500{\rm ms^{-1}}$
$5\mathrm{cm}$
$10\mathrm{cm}$
$5\mathrm{cm}$
$10\mu{ m m}$







Rewriting the phase: $\Phi(z,y) = \tilde{\Phi}(z - \varphi_c(y))$

Center of focal spot in the image





Asymptotic expansion of travel times in the paraxial regime

$$\frac{|z-x_r|}{c} - \frac{|y-x_r|}{c^{\star}} = \sum_k \eta^k f_k(z-\varphi_c(y))$$



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$$\varphi_c(y) = \left(\left(\frac{c}{c^\star}\right)^2 \eta^{\frac{1}{2}} y^\perp, \frac{c}{c^\star} y \right)$$

$$F^{c}(z,y) = \eta^{-1} \left(\frac{c}{c^{\star}}\right)^{2} \frac{\ell_{0}^{4}}{\left(16\pi^{2}|z||\varphi_{c}(y)|\right)^{2}} \int_{\mathcal{B}_{0}} \omega^{2} e^{i\frac{2\omega}{\eta c}(|z|-|\varphi_{c}(y)|)} e^{i\frac{\omega}{c}\left(1-\left(\frac{c^{\star}}{c}\right)^{2}\right)\frac{|\varphi_{c}(y)^{\perp}|^{2}}{|\varphi_{c}(y)|}} \mathcal{G}^{2}\left(\frac{\omega\ell_{0}}{c}\left(\frac{z^{\perp}}{|z|}-\frac{\varphi_{c}(y)^{\perp}}{|\varphi_{c}(y)|}\right), \frac{\omega\ell_{0}^{2}}{c}\left(\frac{1}{|z|}-\left(\frac{c}{c^{\star}}\right)^{2}\frac{1}{|\varphi_{c}(y)|}\right)\right) \mathrm{d}\omega + \mathcal{O}(1).$$

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Analytic approximation of the PSF

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Definition : Confocal \mathcal{F} trace at y_0

$$\begin{split} [c_{\min},c_{\max}] &\longrightarrow \mathbb{C}\\ c &\longmapsto \mathcal{F}(c) := F^{\frac{c}{c^{\star}}}\left(\phi_c(y_0),y_0\right)\\ \end{split}$$
 with $y_0 \in D$

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$$c^{\star} = \operatorname{argmax}_{c} |\mathcal{F}(c)|.$$



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Asymptotic expression

Lemma : Asymptotic development (paraxial regime)

$$\mathcal{F}(c) \sim \mathcal{G}(\beta)^2 \quad \text{with} \quad \beta = \frac{\ell^2 \omega_0}{c^* |y_0|} \left(\left(\frac{c}{c^*}\right)^2 - 1 \right)$$
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Proposition : Phase jump estimator

$$c^{\star} = \operatorname{argmax}_{c} \frac{\partial}{\partial_{c}} \operatorname{Im} \mathcal{F}(c).$$



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Maximising brightness at center of focal spot

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in the medium (*guide star*)

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How does one access the PSF in a microstructured medium

• The PSF can be accessed on the image if there is a strong isolated small reflector

Image domain

 $D' \longrightarrow D'$ $z \longmapsto \left| F^{\frac{c}{c^{\star}}}(z, y_0) \right|$



The dual focal spot

Exchanging the roles of c and c^{\star}

Physical domain

$$D \longrightarrow D$$
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The part of the image affected by the point $y_0 \in D$

Physical domain

$$D \longrightarrow D$$
$$y \longmapsto \left| F^{\frac{c}{c^{\star}}}(z_0, q) \right|$$

Has a small « quasi-support » centered around $y_0(c) = \varphi_c^{-1}(z_0) \in D$



The part of the domain probed by the pixel $z_0 \in D'$

Exchanging the roles of c and c^{\star}

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Virtual random reflector

$$I^{c}(z) = \int_{D} \left(n(y) - n^{\star} \right) F^{\frac{c}{c^{\star}}}(z, y) \mathrm{d}y, \qquad z \in D'$$

 $I^{c}(z) \approx \int_{D(c,z)} \left(n(y) - n^{\star} \right) F^{\frac{c}{c^{\star}}}(z,y) \mathrm{d}y, \qquad z \in D'$



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Proposition : Variance of I^c

$$\mathbb{E}\left[|I^{c}(z)|^{2}\right] = \varepsilon^{d}C_{n}\int_{D(c,z)}|F^{c}(z,y)|^{2}\,\mathrm{d}y + o\left(\varepsilon^{d}\right)$$

$$C_{n} \text{ the integral of the covariance of } n.$$

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 $I^{c}(z)$ depends on the realization of $n(y) - n^{\star}$ in the small subdomain $D(c, z) \subset D$

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$$\mathbb{E}\left[|I^{c}(z)|^{2}\right] = \varepsilon^{d}C_{n}\int_{D(c,z)}|F^{c}(z,y)|^{2}\,\mathrm{d}y + o\left(\varepsilon^{d}\right)$$

$$C_{n} \text{ the integral of the covariance of } n.$$



D(c,z) centered around $\varphi_c^{-1}(z)$ - moves with c

$$I^{c}(z) \approx \int_{D(c,z)} (n(y) - n^{\star}) F^{\frac{c}{c^{\star}}}(z,y) dy,$$

Small domain around $\varphi_{c}^{-1}(z)$



 $D'(\text{Image}) \qquad c = 1.2 * c^* D (\text{Physical domain})$





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 $1540 \text{m/s}, [2,4] \text{MHz}, \star = (0,30) \in (-20,20) \times (0,60) \text{mm}}$



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1 realization

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Pick $t_0 \gg \omega_0^{-1}$ and consider $y_{t_0} := (0, c^* t_0)$ \leftarrow UNKNOWN point in the domain, only t_0 known

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 $c^{\star} = \operatorname{argmax}_{c} \mathbb{E}\left[I^{c}(z(c(t_{0})))^{2}\right]$ with $z(c) = (0, ct_{0})$

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Dual focal spot centered on y_{t_0}

Only depends on t_0 and c, KNOWN



 (ε^d)

Point spread function

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 ε^{d}



Only one realization of the medium

Point spread function



Denote by $n(\cdot, \gamma), \gamma \in \Omega$ the dependency on the randomness of the medium. Consider $\varepsilon \ll |\Delta z| \ll \frac{c}{\omega_0} \sim \frac{c^*}{\omega_0}$.

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$$n(\cdot,\gamma)\big|_{D_{\eta}(c,z)} = n(\cdot,\tau_{\Delta z}\gamma)\big|_{D_{\eta}(c,z+\Delta z)}$$

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Slightly shifted dual focal spot of z

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Possible in practice to access statistics of I^c w/ one realization



Pick a travel time $t_0 \gg \omega_0^{-1}$.

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Compute the matrix $\mathcal{K}_{ij} := |I^{c_j} \left(z(c_j) + (\Delta z)_i \right)|^2$

$$\ll \frac{c_{min}}{\omega_0}.$$

$$\forall c_j \in [c_{min}, c_{max}]$$
$$\forall (\Delta z)_i \in \mathcal{S}$$

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$$\ll \frac{c_{\min}}{\omega_0}.$$

$$(z(c_j) + (\Delta z)_i) |^2 \qquad \forall c_j \in [c_{\min}, c_{\max}]$$

$$\forall (\Delta z)_i \in S$$

Estimator



Simulations by E. **Parolin** (Alpines, INRIA) 166864 scatterers, $\varepsilon = 38.8 \mu m$, f = 4Mhz.

- 1.0e + 00
- 1.04
- 1.03
- 1.02
- 1.01
- 0.99
- 0.98
- 0.97
- 9.6e-01



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Parameters of the simulation

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1.03	Discretisation: 108 DOFs (P3),
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Spatial averaging for a fixed realization



- 1.0e+00 - 1.04 - 1.03 - 1.02 - 1.01 - 1 - 0.99 - 0.98

 $(\Delta z)_i$

- 0.97 - 9.6e-01

 $\frac{n-n^*}{n^*}$



Spatial averaging for a fixed realization



Spatial averaging for a fixed realization





Comparison with experiments





Comparison with experiments





Comparison with experiments



Modelisation

Provided a model for wave propagation (and back scattering) in a random micro-structured medium

Used this model to describe medical ultrasound imaging measurements

Inverse problem

Analysed the point-spread function when there is a mismatch modeled and actual speed of sound in the medium

Characterized the focal spot

Analysed the image quality indicator introduced by Aubry, Bureau, Fink & Al.

Proved that it is possible to access in situ the effective speed of sound in a micro-structured medium from backscattering measurements



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Take Home messages

It is possible to re introduce sparsity by focusing in the medium The amplitude of the PSF is enough to control *a posteriori* the quality of the focusing (at least the geometrical component)

Summary



Speed of sound mapping

Extend the method to more complex speed of sound model

Perspectives



Speed of sound mapping

Extend the method to more complex speed of sound model

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Layered mediums - refraction effects (ongoing work w. C. Prada) Local variation of speed of sound (on going work w. A. Aubry's team)


Perspectives

Speed of sound mapping

Extend the method to more complex speed of sound model

Microstructure analysis

Add some contrast in the divergence (density)

Possible effective anisotropy

Effects of multiple scattering

Extension to propagation in solids (ongoing work w. C. Prada)

Layered mediums - refraction effects (ongoing work w. C. Prada) Local variation of speed of sound (on going work w. A. Aubry's team)

